

## DETERMINANTS OF LAPLACIANS ON THE SPACE OF CONICAL METRICS ON THE SPHERE

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**ABSTRACT.** On a compact surface with smooth boundary, the determinant of the Laplacian associated to a smooth metric on the surface (with Dirichlet boundary conditions if the boundary is nonempty) is a well-defined isospectral invariant. As a function on the moduli space of such surfaces, it is a smooth function whose boundary behavior in certain cases is well understood; see [OPS and K]. In this paper, we restrict ourselves to a certain class of singular metrics on closed surfaces called conical metrics. We show that the determinant of the associated Laplacian is still well defined and that it is a real analytic function on a suitably restricted subset of the space of conical metrics on the sphere.

### 1. INTRODUCTION

Let  $\Sigma$  denote a compact surface without boundary. Let  $\sigma$  be a smooth metric on  $\Sigma$  and  $\Delta$  be its Laplace-Beltrami operator acting on functions. Let  $0 = \lambda_0 < \lambda_1 \leq \lambda_2 \leq \dots$  denote the eigenvalues of  $\Delta$ . Then for an orthonormal basis of functions  $\{u_j\}$  we have  $\Delta u_j + \lambda_j u_j = 0$ . We recall that the determinant of the Laplacian,  $\det' \Delta$ , is formally defined as

$$\det' \Delta = \prod_{j=1}^{\infty} \lambda_j,$$

and that to give meaning to this product we use the standard zeta regularization. We introduce the zeta function

$$(1.1) \quad Z(s) = \sum_{j=1}^{\infty} \lambda_j^{-s}$$

in terms of which

$$(1.2) \quad \det' \Delta = \exp(-Z'(0)).$$

It follows from the asymptotic distribution of the  $\lambda_j$ 's, given by Weyl's law, that  $Z(s)$  defines an analytic function for  $\Re(s) > 1$ . To study the analytic continuation of  $Z(s)$ , we note that since the gamma function satisfies

$$\Gamma(s)a^{-s} = \int_0^{\infty} e^{-at} t^{s-1} dt,$$

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we have for  $\Re(s) > 1 + \varepsilon$ ,

$$\Gamma(s) \sum_{j=1}^{\infty} \lambda_j^{-s} = \sum_{j=1}^{\infty} \Gamma(s) \lambda_j^{-s}.$$

Thus, for  $\Re(s)$  large we can write

$$(1.3) \quad Z(s) = \frac{1}{\Gamma(s)} \int_0^{\infty} \text{TR} \left( e^{\Delta t} - \frac{1}{A} \right) t^s \frac{dt}{t},$$

with  $A$  being the area of  $(\Sigma, \sigma)$ .

It is easy to see that

$$\int_1^{\infty} \left( \sum_{j=1}^{\infty} e^{-\lambda_j t} \right) t^{s-1} dt$$

can be analytically continued in  $s$ , to all of  $\mathbb{C}$ . Also as is well known, see [MS], the kernel of  $e^{\Delta t}$  has the following expansion as  $t \rightarrow 0$ ;

$$(1.4) \quad \sum_{j=0}^{\infty} e^{-\lambda_j t} u_j^2(x) = \frac{1}{4\pi t} + \frac{K(x)}{12\pi} + O(t),$$

where  $K(x)$  is the Gaussian curvature of  $\Sigma$  at  $x$ . Integrating these local invariants over  $\Sigma$  yields as  $t \rightarrow 0$ ,

$$(1.5) \quad \text{TR}(e^{\Delta t}) = \frac{A}{4\pi t} + \frac{\chi(\Sigma)}{6} + O(t),$$

where  $\chi(\Sigma)$  denotes the Euler characteristic of  $\Sigma$ . It follows that

$$(1.6) \quad Z(s) = \frac{1}{\Gamma(s)} \left\{ \frac{A}{4\pi(s-1)} + \left( \frac{\chi(\Sigma)}{6} - 1 \right) \frac{1}{s} + \text{analytic in } s \right\}$$

in the region  $\Re(s) > -1$ . Hence  $Z$  has an analytic continuation as a meromorphic function, is regular at  $s = 0$  and (1.2) makes sense.

In case  $\Sigma$  has smooth boundary, we subject  $\Delta$  to Dirichlet boundary conditions and use the analogue of (1.5) to show that (1.2) still makes sense. We note that to define  $\det' \Delta$  using (1.2) we needed the discreteness of the spectrum,  $\text{spec}(\Delta)$ , the asymptotic distribution of the eigenvalues and some known results about the small time behavior of the heat kernel.

From here on  $\Sigma$  will denote a closed surface of genus  $p$ . Let  $\{P_1, \dots, P_n\}$  be  $n$  distinct points on  $\Sigma$  with  $n \geq 3$  for  $p = 0$  and  $n \geq 1$  for  $p \geq 1$ . A *conical metric* on  $\Sigma$  with vertices  $\{P_1, \dots, P_n\}$  and exponents  $\{a_1, \dots, a_n\}$  is a flat metric on  $\Sigma \setminus \{P_1, \dots, P_n\}$  such that in the neighborhood of each  $P_i$  there are isothermal coordinates so that the line element takes the form

$$(1.7) \quad ds = |w|^{a_i} |dw|$$

in  $0 < |w| < \delta_i$ , for some  $\delta_i > 0$ , where

$$(1.8) \quad \begin{aligned} \text{(a)} \quad & a_i > -1, \\ \text{(b)} \quad & \sum_{i=1}^n a_i = 2p - 2. \end{aligned}$$

We need condition (1.8(a)) to ensure that the line element (1.7) is integrable on smooth curves through the vertices and condition (1.8(b)) follows the Gauss-Bonnet theorem. We shall see that we can still use (1.2) to define  $\det' \Delta$ , when  $\Delta$  is the Laplacian associated to a conical metric on  $\Sigma$ .

Next we say two conical metrics  $\sigma_1$  on  $\Sigma \setminus \{P_1, \dots, P_n\}$  and  $\sigma_2$  on  $\Sigma \setminus \{Q_1, \dots, Q_n\}$  are equivalent if there is a sense-preserving diffeomorphism  $f$  onto  $\Sigma$  onto itself mapping vertices to vertices for which  $\sigma_2 = f^* \sigma_1$ . The space of conical metrics on  $\Sigma$ ,  $\mathcal{E}_n$ , is defined to be

$$\mathcal{E}_n = \{\text{conical metrics}\} / \text{Diff}^+(\Sigma),$$

with the action of the diffeomorphism group as above. Let  $\mathcal{E}_n^*$  be the subset of  $\mathcal{E}_n$  consisting of those metrics with fixed exponents. We show that, for  $\Sigma = S^2$ ,

**Theorem.**  $\det' \Delta$  is real analytic on  $\mathcal{E}_n^*$ .

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## 2. CONIC LAPLACIAN

Let  $\sigma$  be a conical metric on  $\Sigma$  with vertices  $\{P_1, \dots, P_n\}$  and exponents  $\{a_1, \dots, a_n\}$ . In geodesic polar coordinates, this metric takes the form

$$(2.1) \quad ds^2 = dr^2 + (1 + a_i)^2 r^2 d\theta^2, \quad r = |w|^{1+a_i} / (1 + a_i),$$

near  $P_i$ . This is a special case of a cone-like singularity of a manifold where the metric takes the more general form

$$(2.2) \quad \sigma = dr^2 + r^2 \sigma_N(r), \quad 0 \leq r < \varepsilon,$$

with a smooth family of nonsingular metrics  $\sigma_N(r)$  on a smooth compact manifold  $N^m$  without boundary, referred to as the cross section of the cone, and  $\sigma_N(r)$  constant for small  $r$ . In this special case,  $N$  is the circle  $S_{1+a_i}^1$  of radius  $1 + a_i$ . The spectral geometry of Riemannian spaces  $Y$  with cone-like singularities as such has been studied by Cheeger [C], Nagase [N], and Brüning-Seeley [BS]. Melrose [M] has also studied the spectral geometry of *proper conic* metrics defined by the existence of a defining function  $r$ , on a compact manifold with boundary, in terms of which

$$(2.3) \quad \sigma = r^{2s}((dr/r)^2 + h),$$

where  $h$  is a  $C^\infty$ -symmetric 2-tensor. We note that by setting  $w = re^{i\theta}$ , (1.7) becomes a special case of (2.3) with  $s = 1 + a_i$  and  $h = d\theta^2$ .

Let  $X = \Sigma \setminus \{P_1, \dots, P_n\}$  and  $H(X)$  denote the completion of

$$(2.4) \quad \left\{ u \in C_0^\infty(X) : \|u\|_1^2 = \int_X (|\nabla u|^2 + |u|^2) dV < \infty \right\},$$

in the metric induced by  $\|u\|_1^2 = \|u\|^2 + \|\nabla u\|^2$ . On  $H(X)$  we consider the symmetric bilinear form given by

$$(2.5) \quad q(u, v) = \int_X (\nabla u \cdot \nabla v) dV,$$

for  $u, v \in H(X)$ . Then  $q$  is a semibounded closed quadratic form. This gives  $q$  as the quadratic form of a unique selfadjoint operator  $-\Delta$ , i.e.,

$$(2.6) \quad (-\Delta u, v) = q(u, v),$$

where  $u, v \in H(X)$ .

In the general case, with  $Y$  described as above, one has

**Theorem 2.1** [C, N]. *If  $m = \dim N$  is odd, then the collection of eigenfunctions  $\phi$  such that  $\phi$ ,  $d\phi$ , and  $\Delta\phi \in L^2$  determines a complete orthonormal basis of  $L^2$ . The eigenspaces are finite dimensional and the eigenvalues satisfy  $0 \leq \lambda_1 \leq \lambda_2 \leq \dots \rightarrow \infty$ . Furthermore,*

$$N(\lambda) \sim \frac{\text{vol}(Y)\lambda^{(m+1)/2}}{(4\pi)^{(m+1)/2}\Gamma((m+1)/2+1)}, \quad \text{as } \lambda \uparrow \infty,$$

where  $N(\lambda) = |\{j \geq 1 : \lambda_j \leq \lambda\}|$ .

The proof of this theorem (and the corresponding statement if  $\dim N$  is even) uses as a main technique the method of separation of variables and the solutions of Bessel equations with a regular singularity at  $r = 0$ . In the special case of the conic Laplacian on  $S^2$ , the discreteness of  $\text{spec}(\Delta)$  is pretty straightforward:

**Lemma 2.2.** *Let  $\Delta$  be the Laplacian associated with a conical metric  $\sigma$  on  $S^2$ . Then  $\text{spec}(\Delta)$  is discrete.*

*Proof.* Let  $\sigma$  be a conical metric on  $S^2$  with vertices  $\{P_1, \dots, P_n\}$  and exponents  $\{a_1, \dots, a_n\}$ . Let  $B_i$  be the ball centered at  $P_i$  with radius  $r_i$ . Without loss of generality we assume  $r_i = 1$ . Let  $X = S^2 \setminus \{P_1, \dots, P_n\}$ . To show  $\text{spec}(\Delta)$  is discrete, we need only show

$$(2.7) \quad \mathbf{F}_X = \left\{ u \in C^\infty(X) : \int_X (|\nabla u|^2 + u^2) dV \leq 1 \right\}$$

is precompact in  $L^2(X)$ : see [RS].

We show first that for any real number  $\delta$ , the set

$$\mathbf{F}_{B_i} = \left\{ u \in C^\infty(B_i) : \int_{B_i} (|\nabla u|^2 + u^2) dV \leq \delta, \quad u = 0 \text{ on } \partial B_i \right\}$$

is precompact in  $L^2(B_i)$ . For notational convenience, we drop the subscript  $i$ . On  $B$ , the Laplacian associated with  $\sigma$ , acting on functions, is given by

$$(2.8) \quad \Delta = r^{-2a-2}[r\partial_r + r^2\partial_r^2 + \partial_\theta^2].$$

Consider on  $B$  the eigenvalue problem

$$(2.9) \quad \begin{aligned} \Delta u + \lambda u &= 0, \\ u &= 0 \quad \text{on } \partial B. \end{aligned}$$

Set  $u(r, \theta) = f(r)g(\theta)$  and  $\alpha = a + 1$ . Then (2.9) is equivalent to

$$(2.10) \quad \begin{aligned} r^2(f'' + f'/r + \lambda r^{2\alpha}f)/f &= -g''(\theta)/g(\theta), \\ f(1) &= 0. \end{aligned}$$

The first equation in (2.10) gives  $g''(\theta)/g(\theta) = -m^2$ , where  $m$  is a nonnegative integer. Thus  $g(\theta) = \alpha \cos m\theta + \beta \sin m\theta$  for some  $\alpha, \beta$ . Set  $y = f(r)$ . Then

$$(2.11) \quad y'' + (1/r)y' + (\lambda r^{2\alpha} - m^2/r^2)y = 0.$$

The solution of equation (2.11) is given by

$$(2.12) \quad y = c_1 J_\nu(\sqrt{\lambda}r^{\alpha+1}/(\alpha+1)) + c_2 J_{-\nu}(\sqrt{\lambda}r^{\alpha+1}/(\alpha+1)),$$

where  $c_1$  and  $c_2$  are constants,  $\nu = m/(\alpha + 1)$ , and  $J_\nu$  is the Bessel function of order  $\nu$ . Since  $u \in \mathbf{F}_B$ , we have  $c_2 = 0$ . Thus (2.12) reduces to

$$y = c_1 J_\nu(\sqrt{\lambda} r^{\alpha+1}/(\alpha + 1)).$$

The condition  $f(1) = 0$  implies  $\sqrt{\lambda}/(\alpha + 1)$  is a zero of the function  $J_\nu(x)$ , which has infinitely many real zeroes  $k_{m,j}$  ( $j = 1, 2, 3, \dots$ ). We can then write the eigenfunctions in the form

$$J_\nu(k_{m,j} r^{\alpha+1})(\alpha \cos m\theta + \beta \sin m\theta),$$

where  $\alpha, \beta$  are arbitrary. Furthermore any function in  $C^2(B)$  which vanishes on  $\partial B$ , can be expressed in a convergent series of the form

$$(2.13) \quad \sum_{m,j} a_{m,j} J_\nu(k_{m,j} r^{\alpha+1}) \cos m(\theta - \theta_{m,j}).$$

In particular, if  $\{f_k\} \in \mathbf{F}_B$ , then having normalized the eigenfunctions, we get

$$\sum_{m,j} (\lambda_{m,j}^2 + 1) |a_{m,j}^{(k)}|^2 \leq \delta.$$

Thus for a fixed  $m$ , there is a sequence  $\{k_l\}$  such that  $\{a_{r,j}^{(k_l)}\}$  converges for all  $r \leq m$ . Using the diagonal process, there exists  $\{s_l\}$  such that  $\{a_{m,j}^{(s_l)}\} \rightarrow b_{m,j}$  as  $s_l \rightarrow \infty$ . Set

$$f(r, \theta) = \sum_{m,j} b_{m,j} J_\nu(k_{m,j} r^{\alpha+1}) \cos m(\theta - \theta_{m,j}).$$

Then  $\{f_k^{(s_l)}\} \rightarrow f$  in  $L^2(B)$  and  $\mathbf{F}_B$  is precompact in  $L^2(B)$ .

Now we can show that  $\mathbf{F}_X$  is precompact in  $L^2(X)$ . Let  $\{\phi_i\}$  be a partition of unity subordinate to  $B_1, \dots, B_n, X \setminus \bigcup_{i=1}^n V_i$  where  $V_i$  is a closed subset of  $B_i$  and such that  $|\nabla \phi_i| \leq k$  for all  $i$ . Set  $A_0 = X \setminus \bigcup_{i=1}^n V_i$ ,  $A_i = B_i$  for  $1 \leq i \leq n$ . Then for  $u \in \mathbf{F}_X$ ,  $u = \sum_{i=0}^n \phi_i u_i$  and for every  $0 \leq i \leq n$  we have

$$\begin{aligned} \int_{A_i} (|\nabla(\phi_i u)|^2 + (\phi_i u)^2) dV &= \int_{A_i} (|\nabla u \cdot \phi_i + u \cdot \nabla \phi_i|^2 + \phi_i^2 u_i^2) dV \\ &\leq \int_{A_i} (\phi_i^2 |\nabla u|^2 + 2\phi_i u \cdot |\nabla u| |\nabla \phi_i| + u^2 |\nabla \phi_i|^2 + \phi_i^2 u^2) dV \\ &= \int_{A_i} \phi_i^2 (|\nabla u|^2 + u^2) dV + \int_{A_i} (|\nabla \phi_i|^2 u^2 + 2\phi_i u |\nabla u| |\nabla \phi_i|) dV \\ &\leq \int_{A_i} (|\nabla u|^2 + u^2) dV + \int_{A_i} (k^2 u^2 + k(u^2 + |\nabla u|^2)) dV \\ &\leq 1 + k^2 + k, \end{aligned}$$

whenever  $\int_X (|\nabla u|^2 + u^2) dV \leq 1$ . Thus given  $\{g_n\} \in \mathbf{F}_X$ , one has  $\{\phi_i g_n\} \in \mathbf{F}_{A_i}$  for all  $0 \leq i \leq n$ . Using the first result and a diagonal process we can extract a convergent subsequence  $\{\phi_i g_{n_k}\}$  in  $L^2(X)$ , hence the desired claim.  $\square$

### 3. ASYMPTOTICS OF THE HEAT KERNEL

The authors in [BS] give a scheme for computing the asymptotics of  $\text{TR}(e^{-Lt})$  as  $t \rightarrow 0^+$  for certain singular operators  $L$ , a principal example of which is

the Laplace operator for a manifold with a singularity, where the metric takes the form (2.2). The idea there is to study an appropriate power of the resolvent of  $L$  then pass to the heat kernel by a contour integral. The Hilbert space  $L^2$  on which  $\Delta$  acts is a direct sum of an *interior* part  $H_i$ , consisting of those functions vanishing identically within some distance of the singularity, and a *boundary* part  $H_b$ , consisting of the complementary space. Since the interior part is well understood, the heart of the problem is thus the construction of the boundary parametrix. Let  $W = \{w \in \mathbb{C} : |\arg w| < \pi - \varepsilon\}$ ,  $\varepsilon > 0$ . The main result, in our case, is that for any smooth function  $\phi$  supported sufficiently near the singular points, we have

$$(3.1) \quad \text{TR}[\phi(\Delta + z^2)^{-2}] = \int_0^\infty \sigma(x, xz) dx,$$

where  $\sigma(x, w)$  defined on  $\mathbb{R} \times W$  is  $C^\infty$  in  $x$ , with derivatives analytic in  $w$ , and such that  $\sigma(x, w) \sim \sum_j \sigma_j(x) w^{\alpha_j}$  gives the following expansion as  $w \rightarrow \infty$  in  $W$ :

$$(3.2) \quad \begin{aligned} \int_0^\infty \sigma(x, xz) dx &\sim \sum_{k \geq 0} z^{-k-1} \int_0^\infty \frac{w^k}{k!} \sigma^{(k)}(0, w) dw \\ &+ \sum \int_0^\infty (zx)^{\alpha_j} \sigma_j(x) dx \\ &+ \sum_{\alpha_j = -1}^{-\infty} z^{\alpha_j} \ln z \frac{\sigma_j^{(-\alpha_j-1)}(0)}{(-\alpha_j-1)!}. \end{aligned}$$

Here  $\{\alpha_j\}$  is a sequence of complex numbers with  $\Re(\alpha_j) \rightarrow -\infty$ , and the functions  $\sigma^{(k)}(x, w) = \partial_x^k \sigma(x, w)$  and  $\sigma_j \in S(\mathbb{R})$  (Schwartz class) are determined by any interior parametrix valid away from the singular point and the divergent integrals are defined by analytic continuation; see [BS].

To give an idea how these asymptotic expansions arise, the Laplacian in the neighborhood of a vertex  $P_i$  is given by

$$(3.3) \quad \Delta = -\partial_r^2 - r^{-1} \partial_r + r^{-2} \Delta_N,$$

where  $\Delta_N$  is the Laplacian on  $N = S_{1+a_i}^1$ . The change of variables  $f \rightarrow \sqrt{r}f$  transforms  $\Delta$  into

$$(3.4) \quad \Delta = -\partial_r^2 + r^{-2}(\Delta_N - 1/4).$$

Thus on  $H_b$ ,  $\Delta$  is given by

$$(3.5) \quad -\partial_r^2 + r^{-2}A, \quad A = \Delta_N - \frac{1}{4},$$

where the operator  $A$  is unbounded on  $L^2(N)$ , with  $A \geq -\frac{1}{4}$ . The singular operator,  $D_a$ , which is the Friedrich's extension of  $D = -\partial_r^2 + r^{-2}a$ ,  $a \geq -\frac{1}{4}$ , has resolvent with kernel

$$(3.6) \quad k_\nu(x, y, z) = (xy)^{1/2} I_\nu(xz) K_\nu(yz), \quad x \leq y,$$

where  $\nu = \sqrt{a + \frac{1}{4}}$ , and  $I_\nu$ ,  $K_\nu$  are Bessel functions: see [Ca]. If  $\phi$  has compact support, then  $\phi(D_a + z^2)^{-1}$  has finite trace given by an integral as in (3.2) where

$$(3.7) \quad \sigma(x, w) = x\phi(x)I_\nu(w)K_\nu(w)$$

has an expansion in terms of  $w^{-1}, w^{-2}, \dots$  as  $w \rightarrow \infty$  in any open sector  $|\arg w| < \pi/2 - \varepsilon$ . A similar expansion for the trace of  $\phi(\Delta + z^2)^{-2}$ , can be obtained by using the kernel of  $(D_a + z^2)^{-2}$ ,  $a \in \text{spec}(\Delta)$ , which is given on the diagonal by

$$(3.8) \quad k_\nu^2(x, x, z) = \left( -\frac{1}{2z} \frac{\partial}{\partial z} \right) (x I_\nu(xz) K_\nu(xz)).$$

Let  $\gamma$  be a cut-off function supported sufficiently close to the singular point  $x = 0$ , such that  $\Delta = \Delta_b$  on  $\text{supp } \gamma$  and  $\gamma(x) \equiv 1$  for small  $x$ . Then, as shown in [BS],  $\text{TR}(\gamma e^{-\Delta t})$  has an asymptotic expansion, as  $t \rightarrow 0$ , in  $t^{(k-3)/2}$ ,  $t^{(-4-a)/2}$ , and  $t^{(-4-a)/2} \log t$ . Here  $a \leq -1$  and  $k \geq 3$  are integers. Going back to (1.2) and (1.6), we observe that as far as  $\det' \Delta$  is concerned the terms of interest are the constant term,  $c_0$ , and the coefficient of  $\log t$ ,  $c_1$ , in this asymptotic expansion. Problems arise when  $c_1 \neq 0$  since then  $Z(s)$  has a first order pole at  $s = 0$ :

$$\int_0^1 t^{s-1} \log t \, dt = \lim_{\varepsilon \rightarrow 0} \int_\varepsilon^1 t^{s-1} \log t \, dt = \lim_{\varepsilon \rightarrow 0} \left[ \frac{-\varepsilon^s \log \varepsilon}{s} - \frac{1}{s^2} \right] = -\frac{1}{s^2}.$$

Set

$$(3.9) \quad \zeta(s) = \sum_{a \in \text{spec}(A)} \nu(a)^{-s},$$

where  $\nu(a) = \sqrt{a + \frac{1}{4}}$ . Let  $B_k$  denote the  $k$ th Bernoulli number, and  $\text{Res}_k f(z_0)$  denote the coefficient of  $(z - z_0)^{-k}$  in the Laurent expansion of a meromorphic function  $f$  at  $z_0$ . Then  $c_0$  and  $c_1$  are given by

**Theorem 3.1** [BS].

(a)

$$c_0 = -\frac{1}{2} \text{Res}_0 \zeta(-1) - \frac{1}{4} \sum_{k \geq 1} (-1)^k \frac{B_k}{k} \text{Res}_1 \zeta(2k-1) \\ + \frac{\Gamma'(-\frac{1}{2})}{8\sqrt{\pi}} \text{Res}_1 \zeta(-1) + \int_0^\infty g_0(x) \, dx,$$

(b)  $c_1 = \frac{1}{4} \text{Res}_1 \zeta(-1)$ , where  $g_0(x)$  satisfies  $\lim_{x \rightarrow 0} x g_0(x) = -\frac{1}{2} \text{Res}_1 \zeta(-1)$ .

For  $N = S_{1+a}^1$  and  $\sigma_N(r) = (1+a)^2 d\theta^2$ , if  $\beta = 1+a$  then the eigenvalues of  $A$  are given by  $\{n^2/\beta^2 - 1/4\}$ ,  $n \in \mathbb{Z}$ , and the zeta function (3.9) is given by

$$(3.10) \quad \zeta(s) = 2 \sum_{n=1}^\infty \left( \frac{n}{\beta} \right)^{-s} = 2\beta^s \zeta_R(s).$$

Since

$$\zeta_R(-1) = -1/12, \quad \text{Res}_1 \zeta_R(1) = 1, \quad \text{and} \quad B_1 = -1/6,$$

Theorem 3.1 yields  $c_1 = 0$ , and

$$(3.11) \quad c_0 = \frac{1}{12\beta} - \frac{\beta}{12} + \int_0^\infty g_0(x) \, dx.$$

Thus we have the regularized interior term  $\int_0^\infty g_0(x) \, dx$ , plus singular terms in the zeta function (3.9). One can easily show that if  $\sigma$  is flat on  $\Sigma_0 \setminus \{P_1, \dots, P_n\}$

then  $\int_0^\infty g_0(x) dx = 0$ . Since  $c_1 = 0$ ,  $Z(s)$  is regular at  $s = 0$  and we can use (1.2) to define  $\det' \Delta$ .

#### 4. REAL ANALYTICITY

Let  $\mathcal{E}_n$  denote the space of conical metrics on the sphere with  $n$  vertices. Let  $\mathcal{E}_n^*$  denote the subspace consisting of those metrics with fixed exponents. Any point in  $\mathcal{E}_n$  determines a smooth conformal structure on  $S^2$ . Since there is only one conformal structure on  $S^2$  determined by the standard round metric  $\sigma_0$ , it follows that each point in  $\mathcal{E}_n$  has a representative of the form

$$(4.1) \quad \sigma = e^{2\phi} \sigma_0,$$

where  $\phi$  is smooth and harmonic with respect to  $\sigma_0$  on  $S^2 \setminus \{P_1, \dots, P_n\}$ , for some  $P_1, \dots, P_n$ . We identify  $S^2$  with  $\overline{\mathbb{C}} = \mathbb{C} \cup (\infty)$  with local coordinates  $z$  on  $\mathbb{C}$  and  $1/z$  near  $\infty$  and with the standard conformal structure. A conical metric  $\sigma$  on  $\overline{\mathbb{C}} \setminus \{\tau_1, \dots, \tau_n\}$ , with  $\tau_i \in \overline{\mathbb{C}}$ , corresponding to (4.1) may be written as  $\sigma = e^{2\phi} |dz|^2$  on  $\mathbb{C} \setminus \{\tau_1, \dots, \tau_n\}$ , and as  $\sigma = e^{2\phi(1/z)} |dz|^2 / |z|^4$  for  $z$  near 0, using the coordinate  $1/z$  near  $\infty$ . We shall show

**Theorem 4.1.**  $\det' \Delta: \mathcal{E}_n^* \rightarrow \mathbb{R}$  is real analytic.

*Proof.* Let us first prove the theorem for the case of four vertices. Let  $\sigma$  be a conical metric on  $\overline{\mathbb{C}}$  with vertices  $\{\tau, \tau_0, \tau_1, \tau_2\}$  and fixed exponents  $\{a, a_0, a_1, a_2\}$ . Assume that  $\tau_i$  is fixed for  $i = 0, 1, 2$  with  $|\tau_i| > 2$  and  $|\tau| < \varepsilon$  for a given  $\varepsilon > 0$ . Rather than work with  $\Delta_\sigma$ , a fixed operator on a varying space, it is more convenient to work with a family of varying operators on a fixed underlying space. To that end, let  $f$  be a  $C^\infty$ -function with  $f = 1$  inside the circle of radius 1 about 0 and  $f = 0$  outside the circle of radius 2 about 0. Let  $\{g_t(z)\}$  be the flow generated by the vector field  $X = f\tau$  with  $0 \leq t \leq 1$ . Then  $\phi_\tau(z) = g_1(z) = z + f\tau$  is a diffeomorphism of  $\overline{\mathbb{C}}$  fixing  $\tau_i$  and mapping the disk of radius 1 about 0 conformally to the disk of radius 1 about  $\tau$  with  $\phi_\tau(0) = \tau$ . Under pullback of metrics, we get a family of metrics  $\{\sigma_\tau\}$  on  $\overline{\mathbb{C}}$  with fixed vertices  $\{0, \tau_0, \tau_1, \tau_2\}$  which is analytic in  $\tau$  and fixed in the neighborhood of a vertex.

Near a vertex  $\tau$  of a conical metric on surface  $\Sigma$ , we have Figure 1.

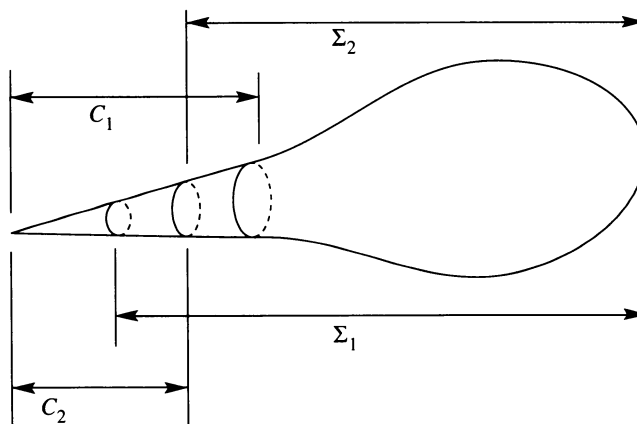


FIGURE 1



It is shown in [N, Lemma 5.3] that

$$(4.2) \quad k_{\tau, \Sigma}(x, x, t) = k_{\tau, \Sigma_1}(x, x, t)I_{\Sigma_2} + k_{\tau, C_1}(x, x, t)I_{C_2} + O(e^{-\delta/t}),$$

where  $k_{\tau, M}$  is the heat kernel on  $M$ ,  $I_M$  is the characteristic function on  $M$  and  $O(e^{-\delta/t})$ ,  $\delta > 0$  is a term any derivatives of which decrease exponentially when  $t \rightarrow 0$ . The main idea there is to reconstruct the heat kernel of  $\Sigma$ , using E. E. Levi's method, from the well-known heat kernel of a compact surface and from the formal representation of the heat kernel involving Bessel functions in a neighborhood of a cone-like singularity. Using estimates of Bessel functions and the uniqueness of the heat kernel, one gets (4.2).

Set  $h(\sigma) = -\log(\det' \Delta)$ , then

$$(4.3) \quad h(\sigma_\tau) = \frac{d}{ds} \Big|_{s=0} \frac{1}{\Gamma(s)} \left[ \int_0^1 \int_{\Sigma} k_{\tau, \Sigma}(x, x, t) dx t^s \frac{dt}{t} \right] + \int_1^\infty \int_{\Sigma} k_{\tau, \Sigma}(x, x, t) dx \frac{dt}{t}.$$

To analyze (4.3) note that we have arranged the problem so that the metric does not change in the neighborhood of a vertex. Thus one parametrix for the Laplacian in a neighborhood of a vertex will work uniformly for all vertices, and the construction of such a parametrix has already been done in [BS]. Away from the vertices, we need a further result. Let  $\{g_\tau\}$  be a family of metrics analytic in  $\tau$  on a compact surface  $M$  with smooth boundary, and  $\{\Delta_\tau\}$  be the corresponding family of Laplacians on  $L^2(M)$ . Let  $\{\phi_{1,\tau}, \phi_{2,\tau}, \dots\}$  be a complete orthonormal basis of  $L^2(M)$  consisting of eigenfunctions of  $\Delta_\tau$  with  $\phi_{j,\tau}$  having eigenvalue  $\lambda_{j,\tau} : 0 < \lambda_{1,\tau} \leq \lambda_{2,\tau} \leq \dots$ . Then  $\lambda_{1,\tau}$  is continuous in  $\tau$  and is uniformly bounded from below by some constant  $c$ . Let  $k_\tau(x, x, t)$  denote the heat kernel on  $(M, g_\tau)$ . Then

**Lemma 4.2.** (1)  $k_\tau(x, x, t)$  is analytic in  $\tau$ .

(2) Given  $T > 0$ ,  $|k_\tau(x, x, t)| = O(e^{-ct/4})$  for all  $t \geq T$ .

(3) For small  $t$  and fixed  $\tau_0$ , the error term in the expansion of  $k_\tau(x, x, t)$  is uniform in  $\tau$  in a neighborhood of  $\tau_0$ .

*Proof.* We recall the construction of a parametrix for  $\Delta$ . Let  $\varepsilon = \text{inj}(M)$ ,  $B_\varepsilon = B(y, \varepsilon)$  for  $y \in M$ . We introduce geodesic spherical coordinates on  $B_\varepsilon$  by

$$x = \exp_y(r\xi), \quad 0 \leq r \leq \varepsilon, \quad \xi \in M_y \text{ with } |\xi| = 1.$$

Let

$$(4.4) \quad S_k(x, y, t) = \frac{1}{4\pi t} \exp\left(\frac{-d^2(x, y)}{4t}\right) \sum_{j=0}^k u_j(x, y) t^j,$$

where  $u_j(\cdot, y) : B_\varepsilon \rightarrow \mathbf{R}$  satisfy

$$(4.5) \quad \begin{aligned} u_0(y, y) &= 1, \\ \partial u_0 / \partial r + \frac{1}{2}(\phi' / \phi) u_0 &= 0, \\ \partial u_j / \partial r + [\frac{1}{2}\phi' / \phi + j/r] u_j &= \Delta u_j / r, \end{aligned}$$

for  $j \geq 1$ . Here we use the notation of [Ch, p. 149], where  $r\phi(\exp_y r\xi, y)$  is defined as a determinant of a path of linear transformations on the orthogonal space of  $\xi$  defined in terms of parallel translation and Jacobi fields along geodesics  $\gamma_\xi$ . The solution to (4.5) is given by

$$(4.6) \quad \begin{aligned} u_0(x, y) &= \phi^{-1/2}(x, y), \\ u_j(x, y) &= u_0(x, y) \int_0^1 \tau^{j-1} (u_0 \Delta u_{j-1})(\exp_y \tau r\xi, y) d\tau, \end{aligned}$$

for  $j \geq 1$ . Let  $0 \leq \rho \leq 1 \in C^\infty(M \times M)$  be equal to 1 on  $B_{\varepsilon/4}$  and 0 on  $(M \times M) \setminus B_{\varepsilon/2}$ . A parametrix for  $(\Delta - \partial/\partial t)$  on  $M$  is then given by  $H_k = \rho S_k$ , which satisfies

$$(4.7) \quad (\Delta_x - \partial/\partial t)H_k = t^{k-1} \exp(-d^2(x, y)/4t)G_k,$$

with  $G_k \in C^\infty(M \times M \times [0, \infty))$ . A fundamental solution to the heat equation is then given by

$$k(x, x, t) = H_k + H_k * \sum_{l=1}^{\infty} \left[ \left( \Delta_x - \frac{\partial}{\partial t} \right) H_k \right]^{*l} = H_k + H_k * F,$$

where for  $0 \leq t \leq T$ , one has the estimate

$$(4.8) \quad |F(x, y, t)| \leq Ct^{k-1} \exp(-d^2(x, y)/4t).$$

It is well known from the theory of ordinary differential equations [CL] that the solution of an ordinary differential equation, depending continuously (or analytically) on a parameter, depends continuously (or analytically) on the given data. Since geodesics  $\gamma: (a, b) \rightarrow M$  are solutions of second degree ordinary differential equations whose coefficients are the Christoffel symbols (which are smooth functions in the coefficients of the metric  $g_\tau$ ), geodesics on  $M$  depend smoothly on  $\tau$ . Similarly for a parallel field of vectors  $v(t)$  along a geodesic or for a Jacobi field  $Z$ . We can thus conclude that  $u_{0,\tau}$  is analytic in  $\tau$ . Using the iteration formula (4.6) we see that  $u_{j,\tau}$  is analytic in  $\tau$  for all  $j$ . This gives (1) in Lemma 4.2.

For the large time behavior, let  $T > 0$  be given. Then

$$(4.9) \quad k_\tau(x, x, t) = \sum_{j=1}^{\infty} e^{-\lambda_{j,\tau} t} \phi_{j,\tau}(x) \phi_{j,\tau}(x),$$

and

$$\begin{aligned} \left| \sum_{j=1}^{\infty} e^{-\lambda_{j,\tau} t} \phi_{j,\tau}(x) \phi_{j,\tau}(x) \right| &\leq \sum_{j=1}^{\infty} e^{-\lambda_{j,\tau} t} \|\phi_{j,\tau}\|_{\infty}^2 \\ &\leq C \sum_{j=1}^{\infty} e^{-\lambda_{j,\tau} t} \|D^2 \phi_{j,\tau}\|_2^2 \leq C \sum_{j=1}^{\infty} e^{-\lambda_{j,\tau} t} (\lambda_{j,\tau})^2, \end{aligned}$$

where we used a standard Sobolev inequality [F]. Using the simple fact that  $e^{-s} s \leq e^{-s/2}$  for  $s \geq \ln 4$ , and the well-known asymptotic distribution of large eigenvalues  $\lambda_j \sim \alpha j$  as  $j \rightarrow \infty$ , we can choose  $N$  large enough so that for  $j \geq N$  we have

$$(4.10) \quad e^{-\lambda_{j,\tau} t} (\lambda_{j,\tau})^2 \leq (2/t^2) e^{-\lambda_{j,\tau} t/4},$$

$$(4.11) \quad \lambda_{j,\tau} \geq \alpha j/2,$$

and

$$(4.12) \quad e^{-(\alpha j/2 - \lambda_{1,\tau})t/4} \leq j^{-2}.$$

Then

$$(4.13) \quad \sum_{j=1}^N e^{-\lambda_{j,\tau}t} (\lambda_{j,\tau})^2 \leq 2N e^{-\lambda_{1,\tau}t/4}$$

and

$$(4.14) \quad \begin{aligned} \sum_{j=N+1}^{\infty} e^{-\lambda_{j,\tau}t} (\lambda_{j,\tau})^2 &\leq 2 \sum_{j=N+1}^{\infty} e^{-\lambda_{j,\tau}t/4} \leq 2e^{-\lambda_{1,\tau}t/4} \sum_{j=N+1}^{\infty} e^{-\mu_{j,\tau}t/4} \\ &\leq 2e^{-\lambda_{1,\tau}t/4} \sum_{j=N+1}^{\infty} e^{-(\alpha j/2 - \lambda_{1,\tau})t/4} \\ &\leq 2e^{-\lambda_{1,\tau}t/4} \sum_{j=N+1}^{\infty} j^{-2} \leq B e^{-\lambda_{1,\tau}t/4}, \end{aligned}$$

where

$$\mu_{j,\tau} = \lambda_{j,\tau} - \lambda_{1,\tau}, \quad B = 2 \sum_{j=N+1}^{\infty} j^{-2}.$$

Putting all this together we get

$$|k_{\tau}(x, x, t)| \leq C(2N + B)e^{-\lambda_{1,\tau}t/4},$$

hence the desired estimate

$$|k_{\tau}(x, x, t)| = O(e^{-ct/4}).$$

For the small time behavior, let  $T > 0$  and  $\tau_0$  be given. Then, for  $t \in [0, T]$ , the constant  $C$  appearing in (4.8) is a function of the volume of  $M$ ,  $T$ , and  $\sup |G_k|$  on  $M \times M \times [0, T]$ . Using the analyticity of  $u_{j,\tau}$  in  $\tau$  on the compact surface  $M$  we get the uniformity in  $\tau$  of  $\sup |G_k|$  in a neighborhood of  $\tau_0$ , hence the desired result. Putting together (4.2), (4.3), Lemma 4.2, and the remark preceding it we get Theorem 4.1 for four vertices.

For the general case, we proceed in a similar manner. Let  $\sigma$  be a conical metric on  $\bar{C}$  with fixed exponents and with vertices  $\{0, 1, \infty, \tau_4, \dots, \tau_n\}$ . Assume that  $|\tau_i - \tau_i^0| < \varepsilon$  for a given  $\varepsilon$  and fixed  $\tau_i^0$ . We construct, as before, a diffeomorphism of  $\bar{C}$  fixing  $\{0, 1, \infty\}$  and mapping a small disk about  $\tau_i$  conformally onto a small disk about  $\tau_i^0$ , with  $\tau_i$  mapped to  $\tau_i^0$ . Under pullback of metrics, we get a family of metrics on  $\bar{C}$  with fixed vertices  $\{0, 1, \infty, \tau_4^0, \dots, \tau_n^0\}$  which is analytic in the parameters  $\tau_i$  and is fixed in the neighborhood of a vertex. The proof is now identical to the one for four vertices. This finishes the proof of the theorem.  $\square$

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